

Griffiths Chapter 1

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Problem 1.3

Consider the gaussian distribution

$$\rho(x) = A \exp(-\lambda(x - a)^2),$$

where A , a , and λ are positive real constants.

- a. Use Equation 1.16 to determine A .

$$1 = \int_{-\infty}^{\infty} A \exp(-\lambda(x - a)^2) dx$$

Pull out constant.

$$1 = A \int_{-\infty}^{\infty} \exp(-\lambda(x - a)^2) dx$$

Let $u = x - a$, $du = dx$

$$1 = A \int_{-\infty}^{\infty} \exp(-\lambda u^2) du$$

From integral table:

$$\int_{-\infty}^{\infty} e^{-cx^2} dx = \sqrt{\frac{\pi}{c}}$$

Substitute:

$$1 = A \sqrt{\frac{\pi}{\lambda}} \implies A = \sqrt{\frac{\lambda}{\pi}}$$

Now we have

$$\rho(x) = \sqrt{\frac{\lambda}{\pi}} \exp(-\lambda(x - a)^2)$$

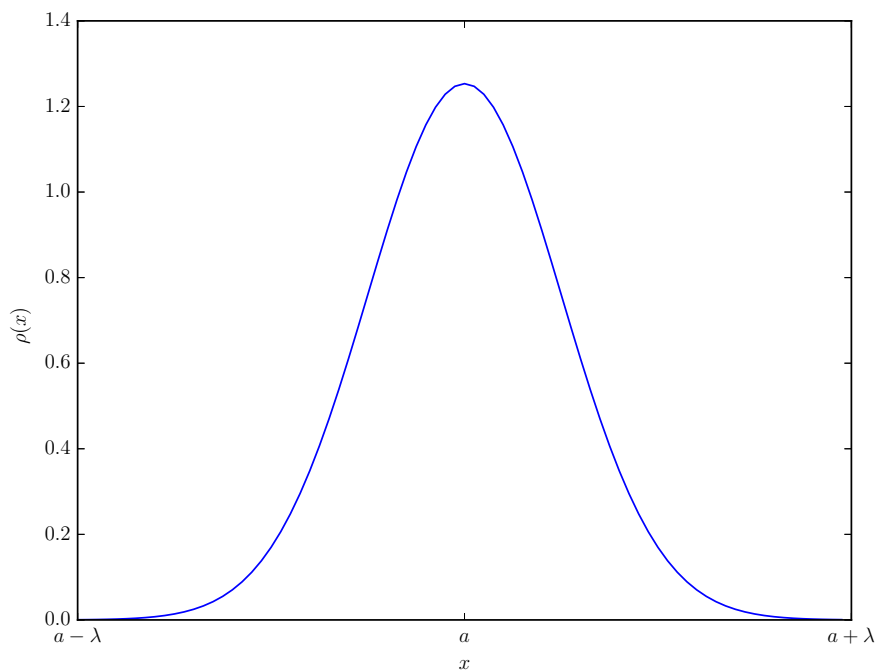
- b. Find $\langle x \rangle$, $\langle x^2 \rangle$, and σ .

$$\langle x \rangle = \int_{-\infty}^{\infty} x \rho(x) dx = \int_{-\infty}^{\infty} x \sqrt{\frac{\lambda}{\pi}} \exp(-\lambda(x - a)^2) dx = \frac{\pi a}{\lambda}$$

$$\langle x^2 \rangle = \int_{-\infty}^{\infty} x^2 \rho(x) dx = \int_{-\infty}^{\infty} x^2 \sqrt{\frac{\lambda}{\pi}} \exp(-\lambda(x - a)^2) dx = \frac{\pi a^2}{\lambda} + \frac{\pi}{2\lambda^2}$$

$$\sigma = \sqrt{\sigma^2} = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \sqrt{\frac{\pi a^2}{\lambda} + \frac{\pi}{2\lambda^2} - \left(\frac{\pi a}{\lambda}\right)^2} = \sqrt{\frac{\pi(2a^2\lambda - 2\pi a^2 + 1)}{2\lambda^2}}$$

- c. Sketch the graph of $\rho(x)$.



Problem 1.4

At time $t = 0$ a particle is represented by the wave function

$$\Psi(x, 0) = \begin{cases} A \frac{x}{a}, & \text{if } 0 \leq x \leq a, \\ A \frac{b-x}{b-a}, & \text{if } a \leq x \leq b, \\ 0, & \text{otherwise,} \end{cases}$$

where A , a , and b are constants.

- a. Normalize Ψ (that is, find A , in terms of a and b).

$$\rho(x, 0) = |\Psi(x, 0)|^2 = \Psi^2(x, 0) = A^2 \begin{cases} \frac{x^2}{a^2}, & \text{if } 0 \leq x \leq a, \\ \frac{(b-x)^2}{(b-a)^2}, & \text{if } a \leq x \leq b, \\ 0, & \text{otherwise} \end{cases}$$

$$1 = \int_{-\infty}^{+\infty} \rho(x) dx = A^2 \int_0^a \frac{x^2}{a^2} dx + A^2 \int_a^b \frac{(b-x)^2}{(b-a)^2} dx$$

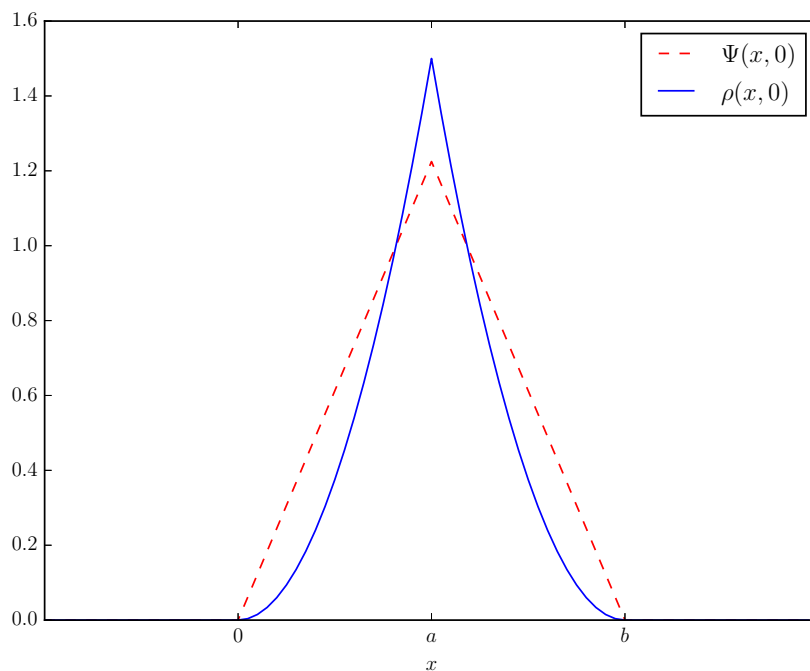
$$u := b-x \implies du = -dx; \quad x = a \implies u = b-a; \quad x = b \implies u = 0$$

$$A^{-2} = \int_0^a \frac{x^2}{a^2} dx - \int_{b-a}^0 \frac{u^2}{(b-a)^2} du = \int_0^a \frac{x^2}{a^2} dx + \int_0^{b-a} \frac{u^2}{(b-a)^2} du$$

$$= \frac{1}{3} \left[\frac{x^3}{a^2} \right]_0^a + \frac{1}{3} \left[\frac{u^3}{(b-a)^2} \right]_0^{b-a} = \frac{1}{3} [a + (b-a)] = \frac{b}{3} \implies A = \sqrt{\frac{3}{b}}$$

$$\Psi(x, 0) = \sqrt{\frac{3}{b}} \begin{cases} \frac{x}{a}, & \text{if } 0 \leq x \leq a, \\ \frac{b-x}{b-a}, & \text{if } a \leq x \leq b, \\ 0, & \text{otherwise,} \end{cases}$$

b. Sketch $\Psi(x, 0)$ as a function of x .



c. Where is the particle most likely to be found, at $t = 0$?

The particle's most likely position is given by $\operatorname{argmax}_x \rho(x, 0)$. To the left of a , ρ is positive and increasing, to the right of it, it is positive and decreasing, and outside the interval $[0, b]$, it is zero, therefore the most likely position is at $x = a$.

- d. What is the probability of finding the particle to the left of a ? Check your result in the limiting cases $b = a$ and $b = 2a$.

$$P[x < a] = \frac{3}{b} \int_0^a \frac{x^2}{a^2} dx = \frac{3}{b} \cdot \frac{1}{3} \left[\frac{x^3}{a^2} \right]_0^a = \frac{a}{b}$$

In the limiting case of $b = a$, this gives a probability of 1, which is to be expected as the probability is 1 over the interval $[0, b]$, which is now the same as $[0, a]$. In the limiting case of $b = 2a$, the probability is $1/2$, which is also expected, as $P(x, 0)$ is symmetric about a when both intervals have equal size, distributing half of the probability on $[0, a]$ and half on $[a, b]$.

- e. What is the expectation value of x ?

$$\begin{aligned} \langle x \rangle &= \int_{-\infty}^{\infty} x \rho(x) dx = \frac{3}{b} \left[\int_0^a \frac{x^3}{a^2} dx + \int_a^b x \frac{(b-x)^2}{(b-a)^2} dx \right] \\ &= \frac{3}{b} \left[\frac{b(2a+b)}{12} \right] = \frac{2a+b}{4} \end{aligned}$$

Problem 1.5

Consider the wave function

$$\Psi(x, t) = A \exp(-\lambda|x|) \exp(-i\omega t),$$

where A , λ , and ω are positive real constants.

- a. Normalize Ψ

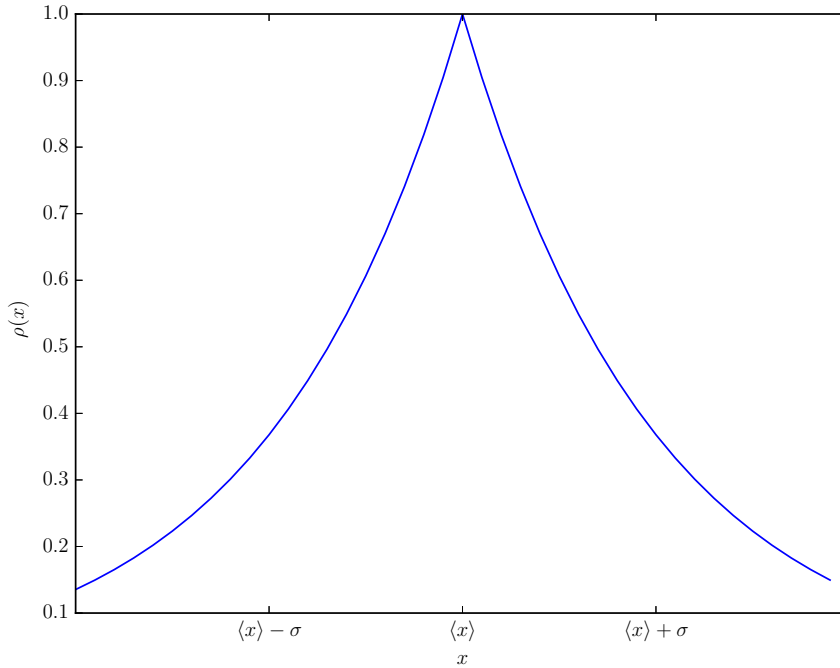
$$\begin{aligned} \rho(x) &= |\Psi(x, t)|^2 = \Psi^* \Psi \\ \Psi &= A \exp(-\lambda|x| - i\omega t) \\ \Psi^* &= A \exp(-\lambda|x| + i\omega t) \\ \rho(x) &= A^2 \exp(-2\lambda|x|) \\ 1 &= \int_{-\infty}^{+\infty} \rho(x) dx \\ A^{-2} &= \int_{-\infty}^{+\infty} \exp(-2\lambda|x|) dx = \frac{1}{\lambda} \\ A &= \sqrt{\lambda} \\ \rho(x) &= \lambda \exp(-2\lambda|x|) \\ \Psi(x, t) &= \sqrt{\lambda} \exp(-\lambda|x| - i\omega t) \end{aligned}$$

- b. Determine the expectation values of x and x^2 .

$$\begin{aligned} \langle x \rangle &= \int_{-\infty}^{+\infty} x \rho(x) dx = \int_{-\infty}^{+\infty} x \lambda \exp(-2\lambda|x|) dx = 0 \\ \langle x^2 \rangle &= \int_{-\infty}^{+\infty} x^2 \rho(x) dx = \int_{-\infty}^{+\infty} x^2 \lambda \exp(-2\lambda|x|) dx = \frac{1}{2\lambda^2} \end{aligned}$$

- c. Find the standard deviation of x . Sketch the graph of $|\Psi|^2$, as a function of x , and mark the points $(\langle x \rangle + \sigma)$ and $(\langle x \rangle - \sigma)$, to illustrate the sense in which σ represents the “spread” in x . What is the probability that the particle would be found outside this range?

$$\sigma = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \sqrt{\langle x^2 \rangle - 0} = \langle x^2 \rangle = \frac{1}{2\lambda^2}$$



The probability of finding the particle *outside* this range is the complement of the probability of finding the particle *inside* this range, which is given by

$$\begin{aligned} \rho[\langle x \rangle - \sigma < x < \langle x \rangle + \sigma]' &= 1 - \rho[\langle x \rangle - \sigma < x < \langle x \rangle + \sigma] \\ &= 1 - \int_{\langle x \rangle - \sigma}^{\langle x \rangle + \sigma} \rho(x) dx \\ &= 1 - \int_{-(2\lambda^2)^{-1}}^{(2\lambda^2)^{-1}} \lambda \exp(-2\lambda|x|) dx \\ &= 1 - (1 - \exp(-\lambda^{-1})) = \exp(-\lambda^{-1}) \end{aligned}$$

Problem 1.7

Calculate $d\langle p \rangle / dt$.

$$\begin{aligned} \langle p \rangle &= -i\hbar \int_{-\infty}^{\infty} \Psi^* \frac{\partial \Psi}{\partial x} dx \\ \frac{d\langle p \rangle}{dt} &= -i\hbar \frac{d}{dt} \int_{-\infty}^{\infty} \Psi^* \frac{\partial \Psi}{\partial x} dx = -i\hbar \int_{-\infty}^{\infty} \left[\Psi^* \frac{\partial}{\partial t} \frac{\partial \Psi}{\partial x} + \frac{\partial \Psi^*}{\partial t} \frac{\partial \Psi}{\partial x} \right] dx \end{aligned}$$

From Schrödinger equation:

$$\begin{aligned}\frac{\partial \Psi}{\partial t} &= \frac{i\hbar}{2m} \frac{\partial^2 \Psi}{\partial x^2} - \frac{i}{\hbar} V \Psi; & \frac{\partial \Psi^*}{\partial t} &= -\frac{i\hbar}{2m} \frac{\partial^2 \Psi^*}{\partial x^2} + \frac{i}{\hbar} V \Psi^* \\ \frac{d\langle p \rangle}{dt} &= -i\hbar \int_{-\infty}^{\infty} \left[\Psi^* \frac{\partial}{\partial t} \frac{\partial \Psi}{\partial x} + \left(-\frac{i\hbar}{2m} \frac{\partial^2 \Psi^*}{\partial x^2} + \frac{i}{\hbar} V \Psi^* \right) \frac{\partial \Psi}{\partial x} \right] dx\end{aligned}$$

Note that the wavefunction has continuous second partial derivatives, and therefore the partial derivatives are *commutative* (by Schwarz' theorem)

$$\begin{aligned}\frac{d\langle p \rangle}{dt} &= -i\hbar \int_{-\infty}^{\infty} \left[\Psi^* \frac{\partial}{\partial x} \frac{\partial \Psi}{\partial t} + \left(-\frac{i\hbar}{2m} \frac{\partial^2 \Psi^*}{\partial x^2} + \frac{i}{\hbar} V \Psi^* \right) \frac{\partial \Psi}{\partial x} \right] dx \\ &= -i\hbar \int_{-\infty}^{\infty} \left[\Psi^* \frac{\partial}{\partial x} \left(\frac{i\hbar}{2m} \frac{\partial^2 \Psi}{\partial x^2} - \frac{i}{\hbar} V \Psi \right) + \left(-\frac{i\hbar}{2m} \frac{\partial^2 \Psi^*}{\partial x^2} + \frac{i}{\hbar} V \Psi^* \right) \frac{\partial \Psi}{\partial x} \right] dx \\ &= -i\hbar \int_{-\infty}^{\infty} \left\{ \Psi^* \left[\frac{i\hbar}{2m} \frac{\partial^3 \Psi}{\partial x^3} - \frac{i}{\hbar} \left(\frac{\partial V}{\partial x} \Psi + V \frac{\partial \Psi}{\partial x} \right) \right] + \left(-\frac{i\hbar}{2m} \frac{\partial^2 \Psi^*}{\partial x^2} + \frac{i}{\hbar} V \Psi^* \right) \frac{\partial \Psi}{\partial x} \right\} dx \\ &= -i\hbar \int_{-\infty}^{\infty} \left\{ \frac{\partial \Psi}{\partial x} \left[\Psi^* \left(-\frac{i}{\hbar} V + \frac{i}{\hbar} V \right) - \frac{i\hbar}{2m} \frac{\partial^2 \Psi^*}{\partial x^2} \right] + \Psi^* \left[\frac{i\hbar}{2m} \frac{\partial^3 \Psi}{\partial x^3} - \frac{i}{\hbar} \frac{\partial V}{\partial x} \Psi \right] \right\} dx \\ &= -i\hbar \int_{-\infty}^{\infty} \left\{ -\frac{i\hbar}{2m} \frac{\partial \Psi}{\partial x} \frac{\partial^2 \Psi^*}{\partial x^2} + \Psi^* \left[\frac{i\hbar}{2m} \frac{\partial^3 \Psi}{\partial x^3} - \frac{i}{\hbar} \frac{\partial V}{\partial x} \Psi \right] \right\} dx \\ &= -i\hbar \int_{-\infty}^{\infty} \Psi^* \left(-\frac{i}{\hbar} \frac{\partial V}{\partial x} \right) \Psi dx - i\hbar \int_{-\infty}^{\infty} \left(\Psi^* \frac{i\hbar}{2m} \frac{\partial^3 \Psi}{\partial x^3} - \frac{i\hbar}{2m} \frac{\partial \Psi}{\partial x} \frac{\partial^2 \Psi^*}{\partial x^2} \right) dx \\ &= \int_{-\infty}^{\infty} \Psi^* \left(-\frac{\partial V}{\partial x} \right) \Psi dx + \frac{\hbar^2}{2m} \left[\int_{-\infty}^{\infty} \Psi^* \frac{\partial^3 \Psi}{\partial x^3} dx - \int_{-\infty}^{\infty} \frac{\partial \Psi}{\partial x} \frac{\partial^2 \Psi^*}{\partial x^2} dx \right]\end{aligned}$$

Let $u = \Psi^*$, then $du = \frac{\partial \Psi^*}{\partial x} dx$. Let $dv = \frac{\partial^3 \Psi}{\partial x^3} dx$, then $v = \frac{\partial^2 \Psi}{\partial x^2}$.

Let $\mu = \frac{\partial \Psi}{\partial x}$, then $d\mu = \frac{\partial^2 \Psi}{\partial x^2} dx$. Let $d\nu = \frac{\partial^2 \Psi^*}{\partial x^2} dx$, then $\nu = \frac{\partial \Psi^*}{\partial x}$.

Using integration by parts, that gives:

$$\begin{aligned}\frac{d\langle p \rangle}{dt} &= \int_{-\infty}^{\infty} \Psi^* \left(-\frac{\partial V}{\partial x} \right) \Psi dx + \frac{\hbar^2}{2m} \left\{ \left(\Psi^* \frac{\partial^2 \Psi}{\partial x^2} \right)_{-\infty}^{\infty} - \int_{-\infty}^{\infty} \frac{\partial^2 \Psi}{\partial x^2} \frac{\partial \Psi^*}{\partial x} dx - \left[\left(\frac{\partial \Psi}{\partial x} \frac{\partial \Psi^*}{\partial x} \right)_{-\infty}^{\infty} - \int_{-\infty}^{\infty} \frac{\partial \Psi^*}{\partial x} \frac{\partial^2 \Psi}{\partial x^2} dx \right] \right\} \\ &= \int_{-\infty}^{\infty} \Psi^* \left(-\frac{\partial V}{\partial x} \right) \Psi dx + \frac{\hbar^2}{2m} \left[\left(\Psi^* \frac{\partial^2 \Psi}{\partial x^2} \right)_{-\infty}^{\infty} - \left(\frac{\partial \Psi}{\partial x} \frac{\partial \Psi^*}{\partial x} \right)_{-\infty}^{\infty} \right]\end{aligned}$$

$\Psi(x, t)$ must go to zero as x goes to (\pm) infinity, so the entire $\hbar^2/2m$ term is zero.

$$\frac{d\langle p \rangle}{dt} = \int_{-\infty}^{\infty} \Psi^* \left(-\frac{\partial V}{\partial x} \right) \Psi dx = \left\langle -\frac{\partial V}{\partial x} \right\rangle$$

We have now arrived at Equation 1.38.