## The Hydrogen Atom

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## **Radial Wave Function**

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# Radial Wave Function in General

- separation of variables rewrites the stationary states of the wave function as  $\psi(\mathbf{r}) = R(r) Y(\theta, \phi)$
- for the sake of simplification, a new variable u is defined, such that u(r) = rR(r)
- the time-independent "radial equation" is thus given by

$$-\frac{\hbar^2}{2m}\frac{\mathrm{d}^2 u}{\mathrm{d}r^2} + \left[V + \frac{\hbar^2}{2m}\frac{\ell(\ell+1)}{r^2}\right]u = Eu.$$

• this is identical to the one-dimensional Schrödinger equation, except the effective potential contains an additional component

$$V_{\text{eff}} = V + \frac{\hbar^2}{2m} \frac{\ell(\ell+1)}{r^2}$$

- the hydrogen atom is comprised by a proton of charge e, and an electron of charge -e
- the proton may be thought to be motionless and centered at the origin
- by Coulomb's law, we may express the potential energy as

$$V(r)=-\frac{e^2}{4\pi\varepsilon_0}\frac{1}{r}$$

• we substitute this into the radial equation

$$-\frac{\hbar^2}{2m}\frac{\mathrm{d}^2 u}{\mathrm{d}r^2} + \left[-\frac{e^2}{4\pi\varepsilon_0}\frac{1}{r} + \frac{\hbar^2}{2m}\frac{\ell(\ell+1)}{r^2}\right]u = Eu.$$

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#### • the Coulomb potential admits both

- continuum states (E > 0), describing electron-proton stattering
- discrete bound states (E < 0), describing the hydrogen atom
- we simplify the notation by introducing  $\kappa$ , defined by

$$\kappa = \frac{\sqrt{-2mE}}{\hbar}$$

• we are only interested in bound states, so E is negative, so  $\kappa$  is real

• dividing the wave equation by E, in terms of m and  $\kappa$ , we obtain

$$\frac{1}{\kappa^2} \frac{\mathrm{d}^2 u}{\mathrm{d}r^2} = \left[ 1 - \frac{me^2}{2\pi\varepsilon_0 \hbar^2 \kappa} \frac{1}{(\kappa r)} + \frac{\ell(\ell+1)}{(\kappa r)^2} \right] u$$

• let  $\rho = \kappa r$  and  $\rho_0 = m e^2 / (2\pi \varepsilon_0 \hbar^2 \kappa)$ 

$$\frac{\mathrm{d}^2 u}{\mathrm{d}\rho^2} = \left[1 - \frac{\rho_0}{\rho} + \frac{\ell(\ell+1)}{\rho^2}\right] u$$

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$$\frac{\mathrm{d}^2 u}{\mathrm{d}\rho^2} = \left[1 - \frac{\rho_0}{\rho} + \frac{\ell(\ell+1)}{\rho^2}\right] u$$

- as  $\rho \to \infty$ , the equation simplifies to

$$\frac{\mathrm{d}^2 u}{\mathrm{d}\rho^2} = u$$

- which has general solution

$$u(\rho) = A \exp(-\rho) + B \exp(\rho)$$

- since  $\exp(\rho) \to \infty$  as  $\rho \to \infty$ , B = 0 for large  $\rho$ 

$$u(\rho) \sim A \exp(-\rho)$$

$$\frac{\mathrm{d}^2 u}{\mathrm{d}\rho^2} = \left[1 - \frac{\rho_0}{\rho} + \frac{\ell(\ell+1)}{\rho^2}\right] u$$

- as  $\rho \to 0$ , the equation simplifies to

$$\frac{\mathrm{d}^2 u}{\mathrm{d}\rho^2} = \frac{\ell(\ell+1)}{\rho^2} u$$

- which has general solution

$$u(\rho) = C\rho^{\ell+1} + D\rho^{-\ell}$$

- as  $\rho \to 0$ ,  $\rho^{-\ell} \to \infty$ , so D = 0

$$u(\rho) \sim C \rho^{\ell+1}$$

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• we introduce a new function,  $v(\rho)$ , defined implicitly by

$$u(\rho) = v(\rho)\rho^{\ell+1}e^{-\rho}$$

- it is simply  $u(\rho)$  stripped of its asymptotic behaviour
- we compute  $\frac{du}{d\rho}$  and  $\frac{d^2u}{d\rho^2}$ , and substitute into the radial equation, giving:

$$\rho \frac{d^2 v}{d\rho^2} + 2(\ell + 1 - \rho) \frac{d v}{d\rho} + [\rho_0 - 2(\ell + 1)]v = 0$$

• the solution to  $v(\rho)$  can be given by the power series

$$v(\rho) = \sum_{j=0}^{\infty} c_j \rho^j$$

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$$v(\rho) = \sum_{j=0}^{\infty} c_j \rho^j$$

- to find the coefficients,  $c_0, c_1, \ldots$ , we first find the derivatives

$$\frac{\mathrm{d}v}{\mathrm{d}\rho} = \sum_{j=0}^{\infty} jc_j \rho^{j-1} = \sum_{j=0}^{\infty} (j+1)c_{j+1}\rho^j$$
$$\frac{\mathrm{d}^2v}{\mathrm{d}\rho^2} = \sum_{j=0}^{\infty} j(j+1)c_{j+1}\rho^{j-1}$$

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• recall the wave equation was given by

$$\rho \frac{d^2 v}{d\rho^2} + 2(\ell + 1 - \rho) \frac{d v}{d\rho} + [\rho_0 - 2(\ell + 1)]v = 0$$

• substituting  $v(\rho)$ ,  $\frac{\mathrm{d}v}{\mathrm{d}\rho}$ , and  $\frac{\mathrm{d}^2v}{\mathrm{d}\rho^2}$  we obtain

$$\sum_{j=0}^{\infty} j(j+1)c_{j+1}\rho^j + 2(\ell+1)\sum_{j=0}^{\infty} (j+1)c_{j+1}\rho^j$$
$$-2\sum_{j=0}^{\infty} jc_j\rho^j + \left[\rho_0 - 2(\ell+1)\right]\sum_{j=0}^{\infty} c_j\rho^j = 0$$

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• dividing through by  $\rho^j$  gives us

$$j(j+1)c_{j+1} + 2(\ell+1)(j+1)c_{j+1} - 2jc_j + [\rho_0 - 2(\ell+1)]c_j = 0$$

• solving for  $c_{j+1}$  gives us the recursive definition

$$c_{j+1} = \frac{2(j+\ell+1) - \rho_0}{(j+1)(j+2\ell+2)} c_j$$

• for large values of j, we have

$$c_{j+1} \approx \frac{2j}{j(j+1)}c_j = \frac{2}{j+1}c_j$$

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$$c_{j+1} \approx \frac{2j}{j(j+1)}c_j = \frac{2}{j+1}c_j$$

- if this approximation were exact, then we would have

$$c_j = \frac{2^j}{j!}c_0$$

- implying

$$v(\rho) = c_0 \sum_{j=0}^{\infty} \frac{2^j}{j!} \rho^j = c_0 e^{2\rho}$$

- so  $u(\rho)$  displays asymptotic behaviour, which we tried to get rid of

- it seems that there is only one way to deal with the issue of asymptotic behaviour: the series must be finite
- there must exist a maximum j, such that

$$c_{j_{\max}+1} = 0$$

• implying

$$2(j_{\max} + \ell + 1) = \rho_0$$

• we now define the **principal quantum number** to be

$$n \equiv j_{\max} + \ell + 1$$

• meaning

$$2n = \rho_0$$

## Results

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### Spectrum of Hydrogen

• energy depends on  $\rho_0$ 

$$E = -\frac{\hbar^2 \kappa^2}{2m} = -\frac{me^4}{8\pi^2 \varepsilon_0^2 \hbar^2 \rho_0^2}$$

• the allowed energies are thus given by

$$E_n = -\left[\frac{m}{2\hbar^2} \left(\frac{e^2}{4\pi\varepsilon_0}\right)^2\right] \frac{1}{n^2} = \frac{E_1}{n^2}$$

where  $n \in \mathbb{Z}^+$ 

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#### Bohr Radius

• recall

$$\rho_0 \equiv \frac{me^2}{2\pi\varepsilon_0\hbar^2\kappa} = 2n$$

• solving for  $\kappa$  gives

$$\kappa = \left(\frac{me^2}{4\pi\varepsilon_0\hbar^2}\right)\frac{1}{n} = \frac{1}{an}$$

• where a is the Bohr Radius

$$a = \frac{4\pi\varepsilon_0\hbar^2}{me^2} \approx 5.29 \times 10^{-11} \,\mathrm{m}$$

• from the definition of  $\rho$  we see

$$\rho = \frac{r}{an}$$

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#### Quantum Numbers

- $\bullet$  we have thus far seen three quantum numbers,  $n,\,\ell,\,{\rm and}\,\,m$
- the spatial wave functions for hydrogen are separated as

$$\psi_{n\ell m}(r,\theta,\phi) = R_{n\ell}(r) Y_{\ell}^{m}(\theta,\phi)$$

• from Section 4.1 we have

$$R_{n\ell}(r) = \frac{1}{r} \rho^{\ell+1} e^{-\rho} v(\rho)$$

• the ground state occurs when n = 1, so the binding energy is given by

$$E_1 = -\left[\frac{m}{2\hbar} \left(\frac{e^2}{4\pi\varepsilon_0}\right)^2\right] = -13.6 \,\mathrm{eV}$$

#### Ground State

• consider  $\psi_{100}(r, \theta, \phi) = R_{10}(r) Y_0^0(\theta, \phi)$ 

$$R_{10}(r) = \frac{c_0}{a} \exp\left(-r/a\right)$$

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 $\bullet$  normalizing gives  $c_0$ 

$$\int_0^\infty |R_{10}|^2 r^2 \, \mathrm{d}r = \frac{|c_0|^2}{a^2} \int_0^\infty \exp(-2r/a) r^2 \, \mathrm{d}r = |c_0|^2 \frac{a}{4} = 1$$
$$c_0 = \frac{2}{\sqrt{a}}$$
$$Y_0^0 = \frac{1}{\sqrt{4\pi}}$$

• the ground state of Hydrogen is given by

$$\psi_{100}(r,\theta,\phi) = \frac{1}{\sqrt{\pi a^3}} \exp(-r/a)$$

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#### Thank You

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